

## HW5: Burgers' equation and the 4/5 Kolmogorov's law

To be returned on February 27, 2018

### I. TRAVELING SHOCK IN BURGERS' EQUATION

During our class, we found the steady shock for the 1D Burgers' equation

$$\partial_t u(x, t) + \frac{1}{2} \partial_x u^2 = \nu \partial_x^2 u, \quad (1)$$

when we maintain the velocities  $u \rightarrow \pm U$  as  $x \rightarrow \mp \infty$ .

1) Generalize the solution to the case when we maintain the velocities at infinity  $u(-\infty) = U_-$  and  $u(+\infty) = U_+$  with  $U_- > U_+$ .

2) Verify that the dissipation remains finite in the limit  $\nu \rightarrow 0$ .

3) The velocity of the shock that you have found above is a special case of the Rankine-Hugoniot condition. By integrating the conservation law

$$\partial_t w(x, t) + \partial_x f(w) = 0, \quad (2)$$

around a shock moving with velocity  $V_s$ , derive the Rankine-Hugoniot relation  $V_s = \frac{f_+ - f_-}{w_+ - w_-}$  where the subscripts  $\pm$  refer to the value on the proximal right and left of the shock, respectively.

### II. PRE-SHOCK IN BURGERS' EQUATION

During our class, we have shown that the first shock occurs at a time  $t^*$  determined by the minimum value of the gradient field at the initial time. The initial gradient and velocity fields are denoted  $g_0(a)$  and  $u_0(a)$ , with  $g_0 = \partial_a u_0$ . Denote the absolute minimum as  $\min_a g_0(a) = -G$  and expand  $g_0(a) \simeq -G + \alpha/2a^2 + \dots$  around its minimum, which is assumed to be at the origin.

1) Denote by  $X(a, t)$  the position at time  $t$  of a particle that was initially at  $a$ . Write down the expression for  $X$  in terms of  $a$  and  $u_0(a)$ .

2) Re-derive the relation between the first shock time  $t^*$  and  $G$ .

3) Show that the inverse Lagrangian map  $a(X, t)$  at  $t = t^*$  behaves singularly around the origin  $X = 0$ , namely it has a 1/3 power-law behavior.

4) Show that at the time  $t^*$  the velocity  $u(X, t)$  also develops a 1/3 singular behavior at the origin.

5) Show that the enstrophy  $\Omega(t) = \int (\partial u / \partial x)^2 dx$  diverges as  $(t^* - t)^{-1/2}$  as  $t$  approaches  $t^*$ .

### III. HOPF-COLE TRANSFORMATION FOR BURGERS' EQUATION

Define the stream function as  $u = -\partial_x \psi(x, t)$ . By taking the space derivative of the Burgers' equation, write down the equation for  $\psi_t$ . An unknown function  $g(t)$  appears when a space-derivative is factored out. Suppose first that the unknown time-dependent function  $g(t) = 0$ .

1) Use the Hopf-Cole transformation  $\psi(x, t) \equiv 2\nu \log \theta(x, t)$  to reduce the previous equation for  $\psi_t$  to the heat equation.

2) Use the Gaussian propagator for the heat equation to obtain the expression of  $\theta(x, t)$ .

3) In the limit  $\nu \rightarrow 0$ , show that  $\psi(x, t) = \max_a \left[ \psi_0(a) - \frac{(x-a)^2}{2t} \right]$ . Interpret the max as follows: considers a parabola  $\frac{(a-x)^2}{2t} + C$  centered at  $x$ , start with a big constant  $C$  and reduce it until you contact the curve  $\psi_0(a)$  for the first time.

4) Interpret double contacts, i.e. first contact at two different  $a$  values for a given  $x$ . What is the velocity profile resulting from multiple  $x$ 's having the same contact point  $a$ ?

5) Can you adapt the arguments in 1) to the case  $g(t) \neq 0$ ?

#### IV. KOLMOGOROV "4/5" LAW FOR BURGERS' EQUATION

By using the same procedure that we discussed in class for the Navier-Stokes equation, derive the relation

$$\langle (u(x) - u(0))^3 \rangle = -12\varepsilon x, \quad (3)$$

for the Burgers' equation. Here,  $\varepsilon$  is the dissipation rate.

#### V. MISSING PARTS IN THE DERIVATION OF THE KOLMOGOROV 4/5 LAW

We shall complete the derivation sketched in class of the 4/5 Kolmogorov law.

1) Use incompressibility to show that the two functions  $B_{NN}(r)$  and  $B_{LL}(r)$  in

$$\langle (v_i(\mathbf{r}) - v_i(\mathbf{0})) (v_j(\mathbf{r}) - v_j(\mathbf{0})) \rangle \equiv B_{NN} \delta_{ij} + (B_{LL} - B_{NN}) \frac{r_i r_j}{r^2}, \quad (4)$$

are related as  $B_{NN} = (1 + r/2\partial_r) B_{LL}$ .

Use this expression to prove that the component proportional to  $\delta_{ij}$  in the time-dependent term  $\partial_t \langle v_i(\mathbf{0}) v_j(\mathbf{r}) \rangle$  and the viscous term  $2\nu \Delta \langle v_i(\mathbf{0}) v_j(\mathbf{r}) \rangle$  in the Navier-Stokes equation read:

$$\left(1 + \frac{r}{2}\partial_r\right) \left(-\frac{2}{3}\varepsilon - \frac{1}{2}\partial_t S_2(r, t)\right); \quad -\nu \left(1 + \frac{r}{2}\partial_r\right) \left(\frac{1}{r^4}\partial_r (r^4 \partial_r S_2(r, t))\right), \quad (5)$$

where the second-order longitudinal structure function  $S_2 = B_{LL}$ .

2) Use incompressibility to show that the three functions  $F_1(r)$ ,  $F_2(r)$  and  $F_3(r)$  in

$$b_{ij,k} \equiv \langle v_i(\mathbf{r}) v_j(\mathbf{r}) v_k(\mathbf{0}) \rangle \equiv F_1 \delta_{ij} \frac{r_k}{r} + F_2 \left(\delta_{ik} \frac{r_j}{r} + \delta_{jk} \frac{r_i}{r}\right) + F_3 \frac{r_i r_j r_k}{r^3}, \quad (6)$$

are related by

$$F_2 = -\left(1 + \frac{r}{2}\partial_r\right) F_1; \quad F_3 = -F_1 + r\partial_r F_1. \quad (7)$$

3) Show that the longitudinal third-order structure  $S_3(r) \equiv \langle (v_i(\mathbf{r}) - v_i(\mathbf{0})) (v_j(\mathbf{r}) - v_j(\mathbf{0})) (v_k(\mathbf{r}) - v_k(\mathbf{0})) \rangle \frac{r_i r_j r_k}{r^3} = 12F_1(r)$ .

4) Show that the component proportional to  $\delta_{ij}$  for the non-linear terms  $\partial_{r_k} [b_{kj,i} - b_{j,ki}]$  in the Navier-Stokes equation reads

$$-2 \left(1 + \frac{r}{2}\partial_r\right) \left(\frac{1}{r^4}\partial_r (r^4 F_1(r, t))\right). \quad (8)$$

5) Gather all bits and pieces to obtain (34.20) in Landau-Lifchitz book

$$-\frac{2}{3}\varepsilon - \frac{1}{2}\partial_t S_2 = \frac{1}{6r^4}\partial_r (r^4 S_3) - \nu \frac{1}{r^4}\partial_r (r^4 \partial_r S_2), \quad (9)$$

which finally yields the 4/5 law.